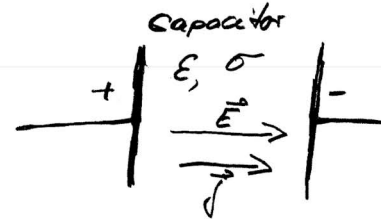


Solutions for Homework # 2

Problem #1: A parallel plate capacitor was charged and then disconnected from the battery. It is discharged due to finite conductivity σ of the media between the plates. The dielectric permittivity is ϵ . Find the conduction current density and the total current density inside the capacitor as a function of time. Find the flux of the electromagnetic energy from the space between the plates. (20 Points)

given are σ and ϵ

a.) find conduction \vec{j}_c and total current \vec{j} !



$$\text{electric field } \vec{E} = \frac{\sigma}{\epsilon} = \frac{Q}{\epsilon \cdot A}$$

$$\text{conduction } \vec{j}_c = \sigma \cdot \vec{E} = \frac{\sigma \cdot Q(t)}{\epsilon \cdot A}$$

$$\text{change of charge } \frac{dQ}{dt} = - \int_S \vec{j} \cdot d^2\vec{r} = -\vec{j} \cdot A = \underline{\underline{-\frac{\sigma \cdot Q}{\epsilon}}}$$

$$\hookrightarrow Q(t) = Q_0 \cdot e^{-t \cdot \sigma / \epsilon} = Q_0 \cdot e^{-t / \tau}; \quad \tau = \epsilon / \sigma$$

$$\hookrightarrow \vec{j}_c = \underline{\underline{\frac{\sigma}{\epsilon \cdot A} \cdot Q_0 \cdot e^{-t / \tau}}}$$

$$\text{diffusion current } \underline{\underline{\vec{j}_D}} = \epsilon_0 \frac{\partial E}{\partial t} = \frac{1}{A} \frac{\partial Q}{\partial t} = \underline{\underline{-\frac{Q_0}{A \cdot \tau} \cdot e^{-t / \tau}}}$$

$$\hookrightarrow \vec{j}_c + \vec{j}_D = 0 = \text{total current}$$

$$\hookrightarrow \vec{B} = 0$$

$$b.) \text{ Flux } \underline{\underline{S}} = \frac{1}{\mu} (\vec{E} \times \vec{B}) = \underline{\underline{0}}$$



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Problem #2: Derivate from the Maxwell equations and the scalar and vector potential definitions the wave propagation equations (20 Points)

Start with Maxwell's equations:

$$\nabla \cdot \vec{D} = \rho, \quad \nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{H} = \vec{j} + \frac{\partial \vec{D}}{\partial t}, \quad \nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

Vector potential

$$\vec{B} = \nabla \times \vec{A} \quad \curvearrowright$$

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} = -\frac{\partial (\nabla \times \vec{A})}{\partial t} \quad \rightarrow \quad \nabla \times \left(\vec{E} + \frac{\partial \vec{A}}{\partial t} \right) = 0$$

$$\hookrightarrow \vec{E} = -\nabla \phi - \frac{\partial \vec{A}}{\partial t}$$

$$\hookrightarrow \text{integrated in } \nabla \times \vec{H} = \vec{j} + \frac{\partial (\epsilon \cdot \vec{E})}{\partial t}$$

$$= \vec{j} + \epsilon \cdot \left(-\nabla \frac{\partial \phi}{\partial t} - \frac{\partial^2 \vec{A}}{\partial t^2} \right) = \vec{j} - \epsilon \nabla \dot{\phi} - \epsilon \ddot{\vec{A}}$$

$$\text{replace } \nabla \times \vec{H} = \frac{1}{\mu} (\nabla \times \vec{B}) = \frac{1}{\mu} (\nabla \times (\nabla \times \vec{A})) = \nabla (\nabla \cdot \vec{A}) - \nabla^2 \vec{A}$$

$$\hookrightarrow \frac{1}{\mu} (\nabla (\nabla \cdot \vec{A}) - \Delta \vec{A}) = \vec{j} - \epsilon \nabla \dot{\phi} - \epsilon \ddot{\vec{A}}$$

$$\rightarrow \boxed{\Delta \vec{A} - \nabla [\nabla \cdot \vec{A} + \mu \epsilon \dot{\phi}] - \mu \cdot \epsilon \ddot{\vec{A}} = -\mu \cdot \vec{j}}$$

Problem #3: Dispersion relations: Problem 7.20 (a) thru (c), Jackson textbook (page 347).
(30 Points)

a.) For a homogeneous, isotropic media with $\mu=1$, a plane wave has to satisfy the Telegraph equation:

$$[\Delta - \mu\sigma - \mu\epsilon] \begin{Bmatrix} \vec{E} \\ \vec{B} \end{Bmatrix} = 0$$

↳ One dim. solution: $u(\omega, n, t) = [A(\omega) e^{\frac{i\omega \cdot n(\omega) \cdot k}{c}} + B(\omega) e^{-\frac{i\omega \cdot n(\omega) \cdot k}{c}}] e^{-i\omega t}$

use Fourier-transform $\omega \rightarrow x$ to get

$$u(x, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-i\omega t} \cdot [A(\omega) e^{i\omega \cdot n \cdot k/c} + B(\omega) e^{-i\omega \cdot n \cdot k/c}] d\omega$$

$$n(\omega) = \sqrt{\frac{c^2}{v^2} - i\mu\sigma c^2/\omega} \quad \text{from Telegraph equ.}$$

b.) $n^*(\omega) = \sqrt{c^2/v^2 + i\mu\sigma c^2/\omega} = \sqrt{c^2/v^2 - i\mu\sigma c^2/(-\omega)} = \underline{\underline{n(-\omega)}}$

c.) put $x=0$ in $u(x, t)$ in a.)

$$u(0, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} [A(\omega) + B(\omega)] e^{-i\omega t} d\omega$$

$$\frac{\partial u(0, t)}{\partial x} = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} [A(\omega) - B(\omega)] \frac{i\omega \cdot n}{c} \cdot e^{-i\omega t} d\omega$$

$$\hookrightarrow A(\omega) + B(\omega) = \frac{1}{\sqrt{2\pi}} \int u(0, t) e^{i\omega t} dt$$

$$A(\omega) - B(\omega) = \frac{1}{\sqrt{2\pi}} \int \frac{-i\omega \cdot n}{\omega \cdot n} \frac{\partial u(0, t)}{\partial x} e^{i\omega t} dt$$

$$\hookrightarrow A(\omega) = \frac{1}{2\sqrt{2\pi}} \int [u(0, t) - \frac{i\omega \cdot n}{\omega \cdot n} \frac{\partial u(0, t)}{\partial x}] e^{i\omega t} dt$$

$$\text{and } B(\omega) = \frac{1}{2\sqrt{2\pi}} \int [u(0, t) + \frac{i\omega \cdot n}{\omega \cdot n} \frac{\partial u(0, t)}{\partial x}] e^{i\omega t} dt$$

$$\Rightarrow \begin{Bmatrix} A(\omega) \\ B(\omega) \end{Bmatrix} = \frac{1}{2\sqrt{2\pi}} \int_{-\infty}^{\infty} [u(0, t) \mp \frac{i\omega \cdot n}{\omega \cdot n} \frac{\partial u(0, t)}{\partial x}] e^{i\omega t} dt //$$

Problem #4: Kramers-Kronig: Problem 7.22 (a) and (b), Jackson textbook (page 348). (30 Points)

Kramers-Kronig

$$\epsilon_1(\omega) = 1 + \frac{2}{\pi} \mathcal{P} \int_0^{\infty} \frac{\omega' \cdot \epsilon_2(\omega')}{\omega'^2 - \omega^2} d\omega'$$

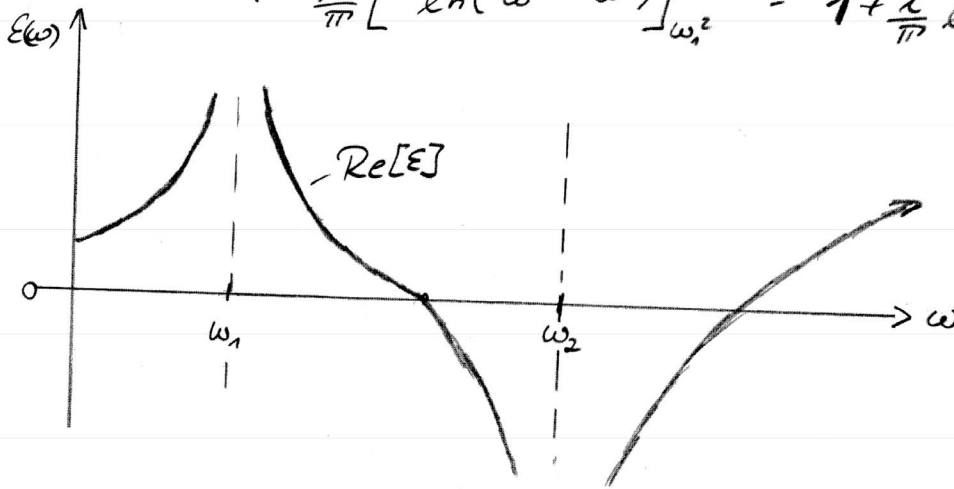
$$\text{and } \epsilon_2(\omega) = -\frac{2\omega}{\pi} \mathcal{P} \int_0^{\infty} \frac{\epsilon_1(\omega') - 1}{\omega'^2 - \omega^2} d\omega'$$

a) $\epsilon_2(\omega) = \lambda \cdot [\Theta(\omega - \omega_1) - \Theta(\omega - \omega_2)]$ with $\omega_1, \omega_2 > 0, \omega_2 > \omega_1$

$$\hookrightarrow \epsilon_1(\omega) = 1 + \frac{2 \cdot \lambda}{\pi} \int_{\omega_1}^{\omega_2} \frac{\omega'}{\omega'^2 - \omega^2} d\omega'$$

$$= 1 + \frac{2 \cdot \lambda}{\pi} \cdot \left(\frac{1}{2} \int_{\omega_1^2}^{\omega_2^2} \frac{d(\omega'^2)}{\omega'^2 - \omega^2} \right)$$

$$= 1 + \frac{\lambda}{\pi} \left[\ln(\omega'^2 - \omega^2) \right]_{\omega_1^2}^{\omega_2^2} = 1 + \frac{\lambda}{\pi} \ln \left(\frac{\omega_2^2 - \omega^2}{\omega_1^2 - \omega^2} \right)$$



b.) $\epsilon_2(\omega) = \frac{\lambda \cdot \gamma \cdot \omega}{(\omega_0^2 - \omega^2)^2 + \gamma^2 \omega^2} \leadsto \epsilon_1(\omega) = 1 + \frac{2}{\pi} \mathcal{P} \int_0^{\infty} \frac{\omega' \cdot \epsilon_2(\omega')}{\omega'^2 - \omega^2} d\omega'$

$$\epsilon_1 = 1 + \frac{2}{\pi} \int_0^{\infty} \frac{\omega' \cdot \lambda \cdot \gamma \cdot \omega}{(\omega'^2 - \omega^2) \cdot [(\omega_0^2 - \omega'^2)^2 + \gamma^2 \omega'^2]} d\omega' \stackrel{\text{Hilbert transformation}}{=} 1 + \frac{\lambda \cdot (\omega_0^2 - \omega^2)}{(\omega_0^2 - \omega^2)^2 + \gamma^2 \omega^2}$$

